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# **Applications of D'alemberts Solution**

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**ABSTRACT:** The method of D'Alembert's provides a solution to the one-dimensional wave equation  $\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$ 

(or)  $u_{xx} = \frac{1}{c^2} u_{tt}$  *i.e.*  $u_{tt} - c^2 u_{xx} = 0$ . In this paper we derive the equation of D'Alembert's solution and also solve the nuclear but the use of D'Alembert's solution

the problem by the use of D'Alembert's solution.

KEYWORDS: Onedimensional wave equation, Cauchy type.

#### I. DERIVATION OF D'ALEMBERT'S SOLUTION

Many authors for example [5] have studied the concept of D'Alemberts method to solve differential equations. The differential equations of homogeneous and non-homogeneous equations[4] Consider the initial value problem of Cauchy type described as[3],

PDE: 
$$u_{tt} - c^2 u_{xx} = 0, -\infty < x < \infty, t \ge 0$$
 (1)

ICs:  $u(x,0) = \eta(x), u_t(x,0) = v(x)$  (2) Where the curve on which the initial data  $\eta(x)$  and v(x) are prescribed as the axis[5].

 $\eta(x)$  and v(x) are assumed to be twice continuously differentiable. Here the string considered as an infinite extent[2]. Let u(x, t) denote the displacement for any x and t.

At time t = 0, the displacement and velocity of the string be prescribed. We know that the general solution of the wave equation is given by[6],

 $u(x,t) = \phi(x+ct) + \psi(x-ct)$ (3)

Where  $\phi$  and  $\psi$  are arbitrary functions. Substituting the ICs (2)in (3), we obtain

$$\phi(x) + \psi(x) = \eta(x)$$
  

$$c[\phi'(x) - \psi'(x)] = v(x)$$

Integrating the second equation of (4), we get

$$\phi(x) - \psi(x) = \frac{1}{c} \int_0^x v(\xi) d\xi$$

Addition and subtraction of this equation with the first relation of Equations (4) yield

$$\phi(x) = \frac{\eta(x)}{2} + \frac{1}{2c} \int_0^x v(\xi) d\xi$$
  

$$\psi(x) = \frac{\eta(x)}{2} - \frac{1}{2c} \int_0^x v(\xi) d\xi$$
 respectively. Substituting these relations for  $\phi(x)$  and  $\psi(x)$  in Eq (2), we ret<sup>[7]</sup>

(4)

at once get[7]

$$u(x,t) = \frac{1}{2} [\eta(x+ct) + \eta(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} v(\xi) d\xi$$
(5)

This is known as the D'Alembert's solution of the one dimensional wave equation.

If v = 0, i.e., if the string is released from the rest, the required solution is[1]



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$$u(x,t) = \frac{1}{2} [\eta(x+ct) + \eta(x-ct)]$$
(6)

The D'Alembert's solution has an interesting geometrical interpretation.

#### II. APPLICATIONS OF D'ALEMBERT'S SOLUTION

#### Example:1

Consider the initial value problem  $u_{tt} = 4u_{xx}$ , t > 0,  $-\infty < x < \infty$  satisfying the conditions u(x, 0) = x,  $u_t(x, 0) = 0[8]$ .

#### Solution:

The D'Alembert's solution for the IVP defind as

PDE:
$$u_{tt} = c^2 u_{xx} t > 0, \quad -\infty < x < \infty$$
  
ICs:  $u(x, 0) = \eta(x)u_t(x, 0) = v(x)$  is  
 $u(x,t) = \frac{1}{2} [\eta(x+ct) + \eta(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} v(\xi) d\xi$ 

For the problem considered, v(x) = 0, c = 2 and  $\eta(x) = x$ . Therefore, the solution to the given problem is

$$u(x,t) = \frac{1}{2}[\eta(x+2t) + \eta(x-2t)] = \frac{1}{2}[x+2t+x-2t] = x$$

Hence the solution of the given IVP isu(x, t) = x.

**Example 2**: Let  $u = \psi(x, t)$  be the solution to the initial value problem

 $u_{tt} = u_{xx}$  for  $-\infty < x < \infty$ , t > 0 with u(x, 0) = sinx,  $u_t(x, 0) = cosx$ , then the value of  $\psi(\frac{\pi}{2}, \frac{\pi}{6})$  is obtained through D'Alembert's solution[9].

Solution:

The D'Alembert's solution for the IVP defind as PDE: $u_{tt} = c^2 u_{xx} t > 0, \quad -\infty < x < \infty$ ICs:  $u(x, 0) = \eta(x)u_t(x, 0) = v(x)$  is  $u(x,t) = \frac{1}{2} [\eta(x+ct) + \eta(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} v(\xi) d\xi$ 

For the given problem, v(x) = cosx, c = 1 and  $\eta(x) = sinx$ .

Therefore, 
$$u(x,t) = \frac{1}{2} [\eta(x+t) + \eta(x-t)] + \frac{1}{2} \int_{x-t}^{x+t} \cos\xi d\xi$$

 $=\frac{1}{2}[\sin(x+t)+\sin(x-t)]+\frac{1}{2}[\sin(x+t)-\sin(x-t)]=\sin 120=\frac{\sqrt{3}}{2}$ 

Hence the solution is  $\sqrt{3}/2$ 

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